

# Constant-Frequency Resonant Inverter for AC-Bus Distribution System

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**Abstract** — A new constant-frequency dc-ac resonant inverter that employs a coupled inductor to achieve ZVS in an entire load and input voltage range with a reduced circulating current is described. The proposed technique is suitable for ac-bus distribution power systems in data-processing equipment. The performance of the proposed approach was evaluated on a 230-W inverter that delivers 1-MHz, 27-V<sub>RMS</sub> ac output voltage from a 320 V - 400 V dc input.

## I. INTRODUCTION

Generally, a high-frequency inverter employs a resonant or a soft-switching technique to reduce switching losses and, consequently, improve conversion efficiency. However, the majority of these resonant-type and soft-switched-type inverters require variable-frequency control to maintain the regulation of their outputs [1]-[3]. While variable-frequency control exhibits acceptable performance in many applications, a number of applications require constant-frequency control. One of the most notable applications that requires a constant-frequency inverter is the ac-bus distribution power system. In such a system, a constant-frequency sinusoidal or trapezoidal ac voltage is distributed to the loads for a final point-of-load conversion by load converters.

Constant-frequency control of inverters is implemented by phase control, which is also called “out-phasing modulation” [4]-[8]. In this type of control, output regulation is achieved at a constant frequency by phase shifting the switching instances of the corresponding switches in the two legs of the inverter. With no phase shift, the output delivers full power, whereas for a phase shift of 180°, the output power is reduced to zero. Generally, the parallel-resonant inverter (PRI) topology and series-parallel resonant inverter (SPRI) topology are employed in applications which require no load operation since the series-resonant inverter (SRI) topology cannot regulate the output at no load.

Figure 1 shows the circuit diagram of a conventional isolated full-bridge PRI. The parallel resonant circuit in Fig. 1 is implemented with primary-side resonant inductor  $L_R$  and secondary-side resonant capacitor  $C_R$ . However, it should be noted that the resonant capacitor could also be placed across the primary winding. Regardless of the placement of the

resonant capacitor, the circuit in Fig. 1 offers a sinusoidal output voltage with relatively low harmonic distortion, as well as output voltage regulation in the entire load range from the full load down to no load. The major deficiency of the PRI is a relatively low partial-load efficiency due to a significant amount of circulating energy in its resonant tank, which is required to maintain output voltage regulation and soft-switching at lighter loads. As a result, the PRI is not suitable for applications that require low power loss at light loads such as power supplies for personal computers.

The light-load performance of the PRI can be improved by employing the SPRI topology shown in Fig. 2 [8]. In the symmetrical SPRI circuit in Fig. 2, which consists of two bridge legs and multiple resonant components, inductors  $L_{R1}$  and  $L_{R2}$  form series resonant circuits with corresponding series capacitors  $C_{S1}$  and  $C_{S2}$  and parallel resonant circuits with capacitor  $C_P$ . Because a properly designed SPRI circuit behaves like a PRI circuit at light loads and like an SRI circuit at higher loads, the SPRI can regulate the output down

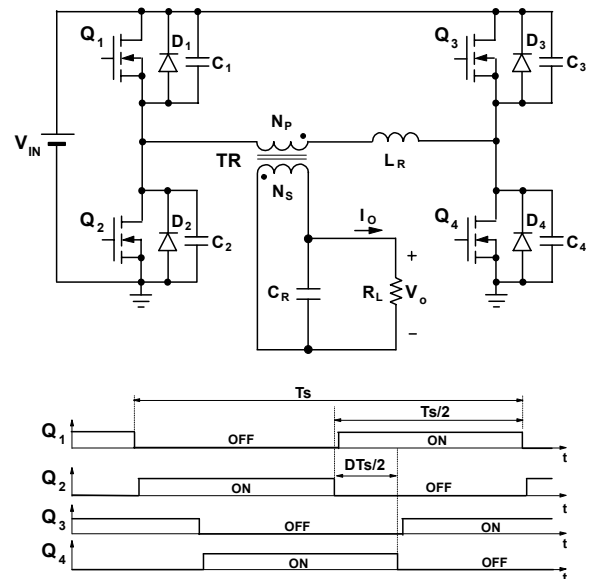


Fig. 1. Conventional parallel resonant inverter and its gate waveforms.

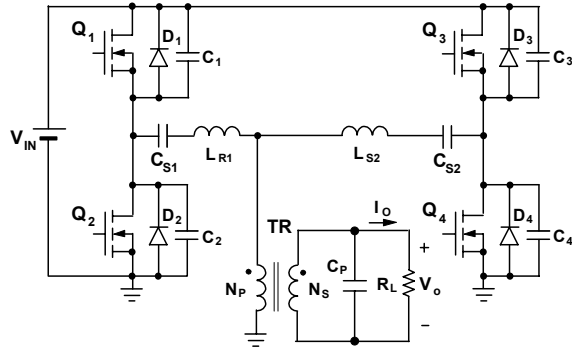


Fig. 2. Conventional class-D isolated phase-controlled series-parallel inverter [8].

to no load with improved partial-load efficiency. Nevertheless, although the circulating energy in the SPRI is reduced compared to that of the PRI, the SPRI still circulates an unnecessarily high amount of energy to maintain soft switching in the entire load and input-voltage range. Moreover, because the Q-factor of a properly designed SPRI that minimizes circulating energy is usually lower than the corresponding Q-factor of the PRI, the output voltage of the SPRI typically exhibits more harmonic distortions than the output voltage of the PRI.

In this paper, a new constant-frequency, dc-ac inverter that employs a coupled inductor to achieve ZVS in a wide range of load currents and input voltages with reduced circulating energy is described. The output voltage regulation in the proposed inverter is achieved by constant-frequency phase-shift control. The performance of the proposed constant-frequency sinusoidal inverter was experimentally verified on a 1-MHz, 230-W prototype that was designed to operate from a 320 V - 400 V dc input and deliver 27-V<sub>RMS</sub> ac output voltage.

## II. PARALLEL RESONANT INVERTER WITH A COUPLED INDUCTOR

The proposed inverter that employs a coupled inductor on its primary side to extend the ZVS range of its switches with minimum circulating energy and conduction loss is shown in Fig. 3. The circuit in Fig. 3 is an isolated PRI, which utilizes a parallel resonant-tank circuit consisting of primary side resonant inductors  $L_{R1}$  and  $L_{R2}$  and secondary-side resonant capacitor  $C_R$  to generate a sinusoidal voltage at the output. One terminal for each of the two resonant inductors is connected to its corresponding inverter leg through a blocking capacitor, while the other terminal is connected with its corresponding winding in coupled inductor  $L_C$ , whose common terminal is connected to transformer TR. The two primary capacitors,  $C_{B1}$  and  $C_{B2}$  are used to prevent saturation in the coupled inductor and in transformer cores by blocking the flow of the dc current through  $L_C$  and TR. Generally, in the PRI these capacitors are selected large

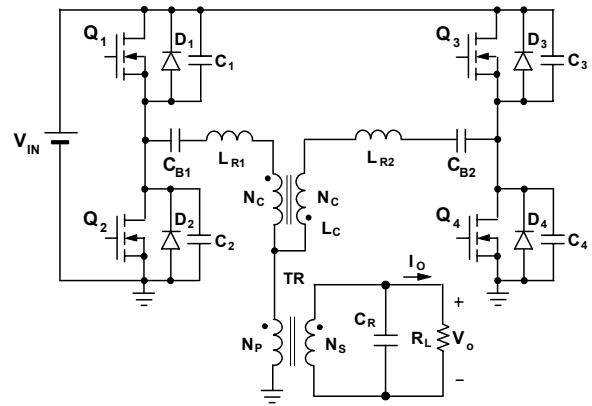


Fig. 3. Proposed parallel resonant inverter with a coupled inductor.

enough so that their voltages are approximately constant during a switching cycle.

The proposed inverter circuit shown in Fig. 3 utilizes the sum of the energy stored in resonant inductors  $L_{R1}$  and  $L_{R2}$  and the magnetizing inductance of coupled inductor  $L_C$  to discharge the capacitance across the switch that is about to be turned on and, consequently, achieve ZVS. Moreover, since coupled inductor  $L_C$  can transfer current (energy) from the winding in one bridge leg to another, the circuit has the unique characteristic that all switches are turned off when they carry a current of the same magnitude. As a result, the energy available for the discharge of the capacitance of the switch is the same for all switches. In addition, the proposed circuit can achieve ZVS of the switches at no load by properly selecting the value of the magnetizing inductance of the coupled inductor. In the proposed circuit, since the energy required for ZVS is mainly provided by the energy stored in the magnetizing inductance of the coupled inductor, the circulating energy of the resonant tank circuit can be

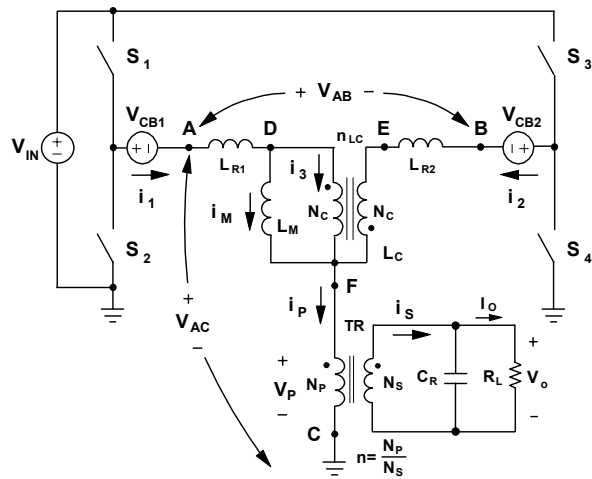


Fig. 4. Simplified circuit diagram of the proposed inverter shown in Fig. 3 along with reference directions of key currents and voltages.

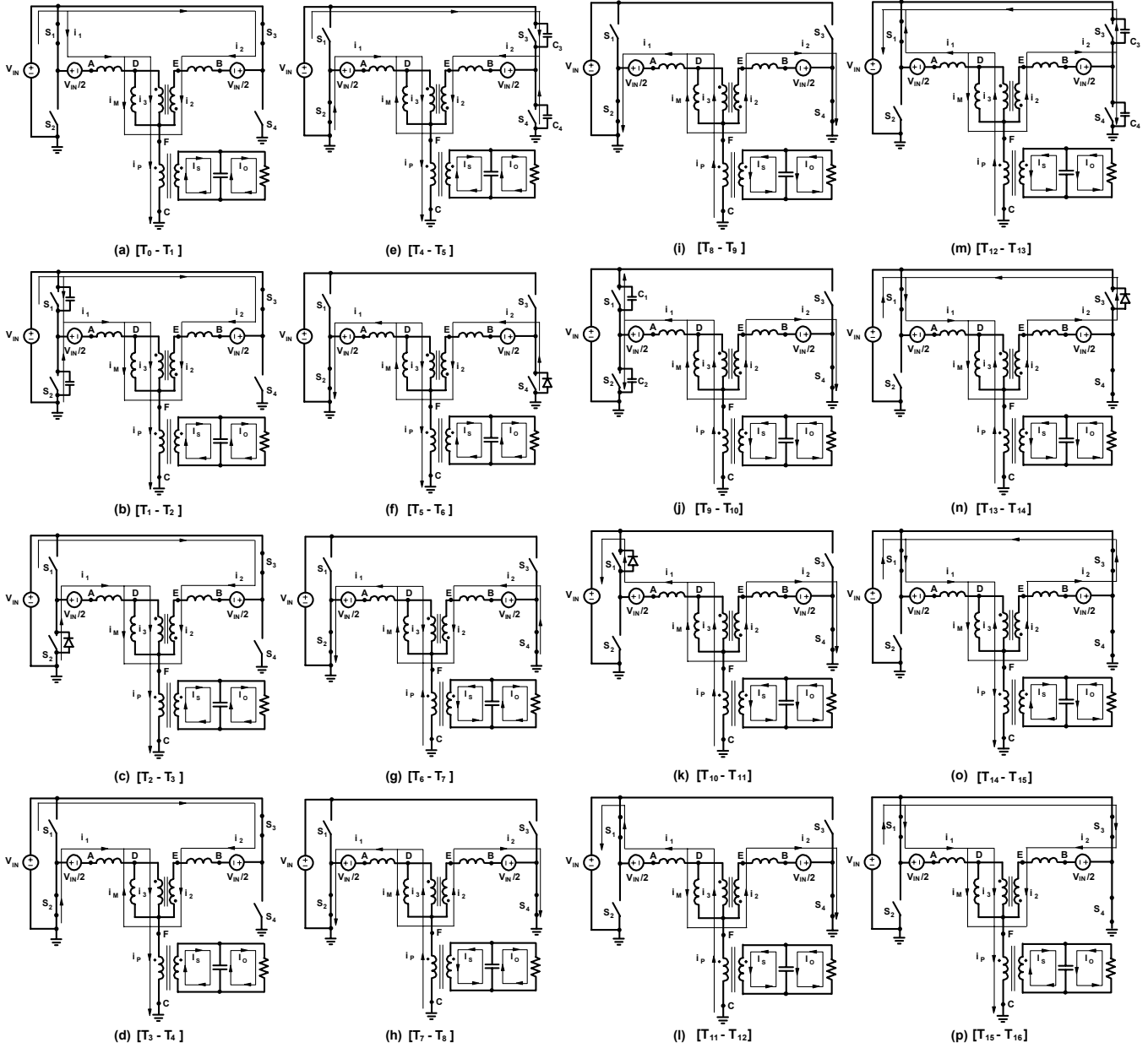


Fig. 5. Topological stages during a switching period.

minimized, which significantly improves partial-load efficiency.

To facilitate the explanation of the operation of the circuit in Fig.3, Fig. 4 shows its simplified circuit diagram. In the simplified circuit, it is assumed that the capacitance of blocking capacitors  $C_{B1}$  and  $C_{B2}$  is large enough so that the capacitors can be modeled as constant voltage sources. Because the average voltages of the coupled inductor windings and the transformer windings during a switching cycle are zero and because the pair of switches in each bridge leg operates with 50% duty cycle when phase-shift control is

used, the magnitude of voltage sources  $V_{CB1}$  and  $V_{CB2}$  in Fig. 4 is equal to  $V_{IN}/2$ , i.e.,  $V_{CB1}=V_{CB2}=V_{IN}/2$ .

To further simplify the analysis, it is also assumed that the resistance of the conducting semiconductor switches is zero, whereas the resistance of the non-conducting switches is infinite. In addition, small leakage inductances in the coupled inductor and transformer, as well as a large magnetizing inductance in the transformer are neglected because their effect on the operation of the converter is negligible. The magnetizing inductance of coupled inductor  $L_C$  and output capacitances of primary switches  $C_1 - C_4$  are not neglected in this analysis since they play a major role in the operation of

the circuit. In Fig. 4, coupled inductor  $L_C$  is modeled as an ideal transformer with turns ratio  $n_{LC}=1$  and with magnetizing inductance  $L_M$  parallel to winding DF. The number of turns for each of the windings of  $L_C$  is  $N_C$ .

Finally, to further facilitate the analysis, Fig. 5 shows the topological stages of the converter during a switching cycle, whereas Fig. 6 shows the key waveforms.

As can be seen from Fig. 4, because the turns ratio of the windings of  $L_C$  is unity ( $n_{LC}=1$ ), current  $i_3$  flowing through winding DF is always equal to current  $i_2$  flowing through winding EF, *i.e.*,  $i_3=i_2$ . As a result,

$$i_1 = i_3 + i_M = i_2 + i_M, \quad (1)$$

and since

$$i_P = i_1 + i_2 \quad (2)$$

the relationship between currents  $i_1$ ,  $i_2$ ,  $i_M$ , and  $i_P$  can be expressed as

$$i_1 = \frac{1}{2}(i_P + i_M) \quad (3)$$

and

$$i_2 = \frac{1}{2}(i_P - i_M). \quad (4)$$

As shown in Fig. 6, during the time interval  $T_0$ – $T_1$ , switch  $S_1$  in leg  $S_1$ – $S_2$  and switch  $S_3$  in leg  $S_3$ – $S_4$  are closed and currents  $i_1$  and  $i_2$  flow through the corresponding switch, blocking capacitor, resonant inductor, and winding of the coupled inductor into the primary of the transformer, as can be seen in the equivalent circuit in Fig. 5(a). From Fig. 5(a), it can also be seen that voltage  $v_{AB}$  must be zero during this topological stage since voltage sources  $V_{CB1}=V_{IN}/2$  and  $V_{CB2}=V_{IN}/2$  are connected in opposition through closed switches  $S_1$  and  $S_3$ . Furthermore, because of the winding orientation of coupled inductor  $L_C$  (dot positions in Fig. 5(a)),

$$\begin{aligned} v_{AB} &= v_{AD} + v_{DF} + v_{FE} + v_{EB} = \\ L_{R1} \frac{di_1}{dt} + v_{DF} + v_{FE} - L_{R2} \frac{di_2}{dt} &= 0 \end{aligned} \quad (5)$$

can only be maintained if the voltages across the couple inductor windings are zero. Namely, for a relatively large magnetizing inductance  $L_M$ , magnetizing current  $i_M$  is virtually constant during the  $T_0$ – $T_1$  interval so that  $di_1/dt=di_2/dt$  and, therefore, for  $L_{R1}=L_{R2}$ ,  $v_{AD}=-v_{EB}$ . As a result, relationship  $v_{AB}=v_{AD}+v_{DF}+v_{FE}+v_{EB}=v_{DF}+v_{FE}=0$  can only be fulfilled if  $v_{DF}=v_{FE}=0$  since both  $v_{DF}$  and  $v_{FE}$  are positive (*i.e.*, have same polarity). As shown in Fig. 6(j), in this topological stage  $v_{AC}=v_{BC}=V_{IN}-V_{IN}/2=V_{IN}/2$ .

When at  $t=T_1$ , switch  $S_1$  is turned off, and current  $i_1$  is diverted from the transistor of switch  $S_1$  to its output capacitance  $C_1$ , as shown in Fig. 5(b). In this topological stage, current  $i_1$  charges capacitor  $C_1$  and discharges capacitor  $C_2$  at the same rate since the sum of the capacitor voltages is equal to constant voltage  $V_{IN}$ , as illustrated in Figs. 6(e) and (f). As a result, the potential of point A starts decreasing causing a decrease of voltages  $v_{AB}$  and  $v_{AC}$ .

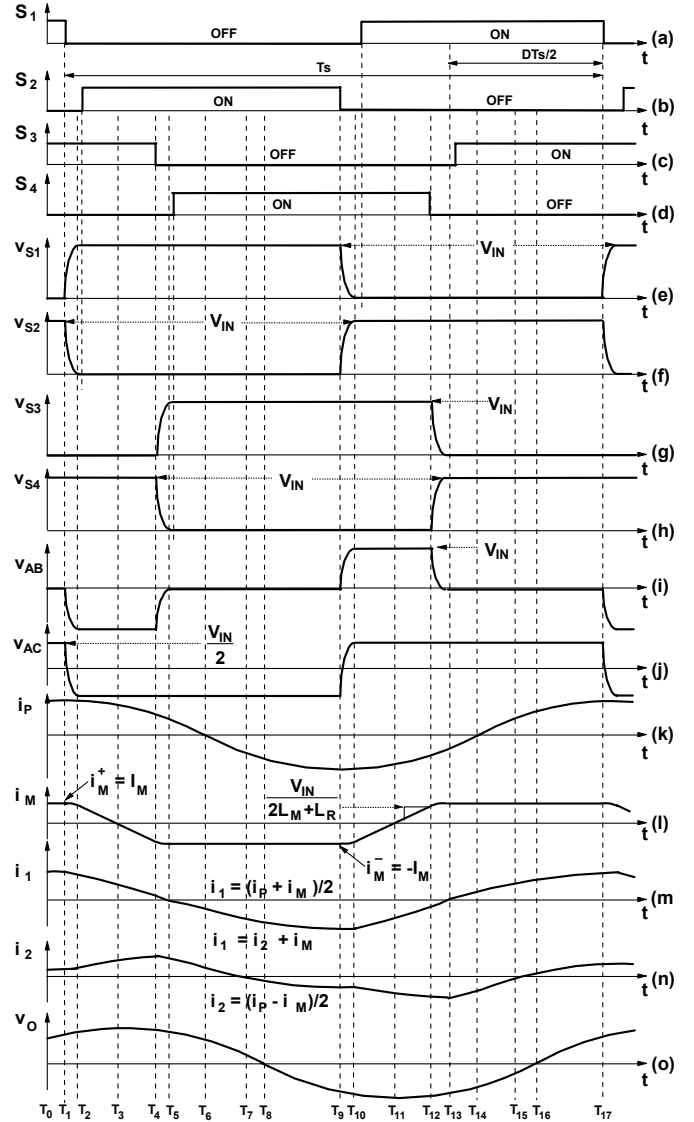


Fig. 6. Key waveforms.

Namely,  $v_{AB}$  decreases from zero toward negative  $V_{IN}$ , whereas  $v_{AC}$  decreases from  $V_{IN}/2$  toward negative  $V_{IN}/2$ , as illustrated in Figs. 6(i) and (j). After capacitor  $C_2$  is fully discharged, *i.e.*, when voltage  $v_{S2}$  reaches zero, current  $i_1$  begins flowing through antiparallel diode  $D_2$  of switch  $S_2$ , as shown in Fig. 5(c). Due to negative voltage  $V_{IN}/2$  applied across series connected resonant inductor  $L_{R1}$  and magnetizing inductance  $L_M$  of coupled inductor  $L_C$ , magnetizing current  $i_M$  decreases with a rate of approximately  $(V_{IN}/2)/(L_M+L_{R1})$ . At the same time, primary current  $i_P$  decreases in a resonant fashion. As a result, current  $i_1=(i_P+i_M)/2$  decreases, whereas current  $i_2=(i_P-i_M)/2$  increases. To achieve zero-voltage turn on of  $S_2$ , it is necessary to turn on  $S_2$  while its antiparallel diode  $D_2$  is conducting. In Fig. 6,  $S_2$  is turned on immediately after  $v_{S2}$  has fallen to zero.

After  $t=T_3$ , when magnetizing current  $i_M$  becomes zero, its magnitude continues to increase in the negative direction, as shown in Fig. 5(d). As a result, current  $i_1$  continues to decrease, whereas current  $i_2$  continues to increase, as seen from waveforms in Figs. 6(m) and (n). At  $t=T_4$ , switch  $S_3$  is turned off so that current  $i_2$  is diverted from the transistor of switch  $S_3$  to its output capacitance  $C_3$ , as shown in Fig. 5(e). During this transition  $C_3$  is charging, while  $C_4$  is discharging at the same rate, voltage  $v_{S3}$  increases from zero toward  $V_{IN}$ , whereas voltage  $v_{S4}$  decreases from  $V_{IN}$  to zero, as illustrated in Figs. 6(g) and (h). Also, since, during this topological stage, the potential of point B decreases from  $V_{IN}/2$  to  $-V_{IN}/2$  while the potential of point A is constant at  $-V_{IN}/2$ , voltage  $v_{AB}$  increases from  $-V_{IN}$  to zero. To achieve ZVS of switch  $S_4$ , the switch needs to be turned on while its antiparallel diode is conducting, *i.e.*, during the  $T_5$ - $T_6$  interval. In Fig. 6,  $S_4$  is turned on immediately after  $t=T_5$ , *i.e.*, immediately after  $v_{S4}$  falls to zero to minimize the conduction loss since the switch conduction loss is lower than the conduction loss of its antiparallel diode. At  $t=T_6$ , resonant primary current  $i_p$  changes direction from positive to negative, and the inverter enters the topological stage shown in Fig. 5(g). Because of resonant capacitor  $C_R$  connected in parallel with the load, the output voltage lags the primary current, *i.e.*, it is still positive at  $t=T_6$ . During the stage in Fig. 5(g), currents  $i_1$  and  $i_2$  continue to decrease. After current  $i_2$  reaches zero at  $t=T_7$ , it continues to increase in the negative direction, as shown in Fig. 6(n). This topological stage, shown in Fig. 5(h), ends at  $t=T_8$  when sinusoidal output voltage  $v_O$  crosses zero. During the next topological stage, shown in Fig. 5(i), the output voltage continues to resonate in the negative direction.

The second half of a switching cycle starts at  $t=T_9$ , when  $S_2$  is turned off. If the switching frequency of the inverter is higher than the resonant frequency of the resonant-tank circuits, switch  $S_2$  is turned off when currents  $i_1$ ,  $i_2$ , and  $i_p$  are negative. As a result, after  $S_2$  is turned off, negative current  $i_1$  charges capacitance  $C_2$  of switch  $S_2$  and discharges capacitance  $C_1$  of switch  $S_1$ , as shown in Fig. 5(j). During this switching transition, voltage  $v_{AB}$  increases from zero to  $V_{IN}$ , while primary voltage  $v_{AC}$  increases from  $-V_{IN}/2$  to  $V_{IN}/2$ . This topological stage ends at  $t=T_{10}$  when voltage  $v_{S1}$  across switch  $S_1$  reaches zero and antiparallel diode  $D_1$  of switch  $S_1$  starts conducting current  $i_1$ , as shown in Fig. 5(k). To achieve ZVS of switch  $S_1$ , switch  $S_1$  needs to be turned on while  $D_1$  is conducting. In Fig. 6, switch  $S_1$  is turned on immediately after  $v_{S1}$  has fallen to zero. Since, after switch  $S_2$  is turned off, voltage  $v_{AB}$  starts increasing, magnetizing current  $i_M$  starts increasing as well, as can be seen from Fig. 6(l). At  $t=T_{11}$ ,  $i_M$  becomes positive, as shown in both Fig. 5(l) and Fig. 6(l). Finally, at  $t=T_{12}$  switch  $S_4$  is turned off, which initiates switching transition in the  $S_3$ - $S_4$  leg. Because, during this transition,  $C_3$  is discharging and  $C_4$  is charging, the potential of point B is increasing from  $-V_{IN}/2$  to  $V_{IN}/2$ . Since, during this time, the potential of point A is constant at  $V_{IN}/2$ , voltage  $v_{AB}$  is decreasing from  $V_{IN}$  to zero. At  $t=T_{13}$

capacitance of switch  $S_3$  is fully discharged and current  $i_2$  begins flowing through antiparallel diode  $D_3$  of switch  $S_3$ , as shown in Fig. 5(n). To achieve ZVS, switch  $S_3$  is turned on shortly after  $D_3$  starts conducting. During the topological stage in Fig. 5(n), primary current  $i_p$ , current  $i_1$ , and current  $i_2$  continue to increase from negative values toward positive values, as seen from waveforms in Figs. 6(k), (m), and (n). Current  $i_1$  becomes positive at  $t=T_{14}$ , whereas current  $i_2$  becomes positive at  $t=T_{15}$ , which completes a cycle of operation of the circuit. The next switching cycle is initiated at  $t=T_{17}$ .

As can be seen from the waveforms in Fig. 6, the commutation of the switches in the  $S_1$ - $S_2$  leg is initiated when current  $i_1 = i_2 + i_M = (i_p + i_M)/2$  is maximal, *i.e.*, when  $i_1 = (i_p + I_M^+)/2$ . Also, the commutation of the switches in the  $S_3$ - $S_4$  leg is initiated when current  $i_2 = (i_p - i_M)/2$  is maximal, *i.e.*, when  $i_2 = (i_p - I_M^-)/2 = (i_p + I_M^+)/2$ . Therefore, in the proposed circuit, all primary switches are commutated with approximately the same magnitude current, which is a unique characteristic of the circuit. Specifically, the charging and discharging of the capacitances of switches in the  $S_1$ - $S_2$  leg is done by the sum of the energy stored in resonant inductor  $L_{R1}$  and the energy stored in the magnetizing inductance of coupled inductor  $L_C$ , whereas the charging and discharging of the capacitances of switches in the  $S_3$ - $S_4$  leg is done by the sum of the energy stored in  $L_{R2}$  and  $L_M$ . To achieve ZVS of the bridge switches, the sum of the energy stored in the magnetizing inductance of the coupled inductor and the energy stored in the resonant inductor of the corresponding bridge leg must be high enough to fully discharge the capacitance of the switch that is about to be turned on. Generally, this ZVS condition can be expressed as

$$\frac{1}{2} \left[ \frac{1}{2} L_R i_p^2 + \frac{1}{2} L_M I_M^2 \right] \geq C V_{IN}^2, \quad (6)$$

where  $C=C_1=C_2=C_3=C_4$  is the total capacitance across the primary switch, which, beside output capacitance  $C_{OSS}$  of the switch, also includes any externally added capacitance.

Since from Fig. 4, primary current  $i_p$  is

$$\vec{i}_p = \frac{1}{n} \cdot \vec{v}_O \left( \frac{1}{R_L} + \frac{1}{j \cdot 2 \cdot \pi \cdot f_S \cdot C_R} \right), \quad (7)$$

it can be seen that energy stored in resonant inductors  $L_{R1}$  and  $L_{R2}$  is dependent on the load, as well as resonant capacitor  $C_R$ . At no load ( $R_L = \infty$ ), the primary current is the reflected current flowing through resonant capacitor  $C_R$ . This current represents the circulating current of the converter that is used to provide the energy required for the resonant inductors to achieve ZVS. In the proposed circuit, this circulating current can be minimized because all of the ZVS energy can be provided by the energy stored in the magnetizing inductance of the transformer. If the energy

stored in the magnetizing inductance is selected so that at the maximum input voltage  $V_{IN(max)}$ ,

$$L_M I_M^2 \geq 4CV_{IN(max)}^2, \quad (8)$$

ZVS will be achieved in the entire load and input-voltage range regardless of the value of  $C_R$ .

The value of  $L_M$  required to achieve ZVS at no load can be calculated from Fig. 6 by observing that, during time interval  $T_{10}$ - $T_{12}$ , magnetizing current  $i_M$  changes linearly from maximum negative value  $-I_M$  to maximum positive value  $I_M$ , *i.e.*,  $i_M$  changes for  $2I_M$ , due to a positive voltage of  $V_{IN}/2$  impressed across the series connection of  $L_{R1}$  and  $L_C$ . From the equivalent circuit in Fig. 5(k), it can be derived that the rate of change of magnetizing current  $i_M$  during the  $T_{10}$ - $T_{12}$  time interval is

$$\frac{di_M}{dt} = \frac{V_{IN}}{2L_M + L_R}, \quad (9)$$

where  $L_R=L_{R1}=L_{R2}$ .

Since, according to Fig. 6, the time interval  $T_{10}$ - $T_{12}$  is approximately equal to

$(1-D)T_S/2$ , where  $D$  is a duty cycle and  $T_S$  is a switching period,  $I_M$  can be calculated from

$$V_{IN} = (2L_M + L_R) \cdot \frac{2I_M}{(1-D) \frac{T_S}{2}} \quad (10)$$

as

$$I_M = \frac{(1-D)V_{IN}}{4 \cdot (2L_M + L_R) \cdot f_S}, \quad (11)$$

where  $f_S=1/T_S$  is the switching frequency. Since  $D \approx 0$  at no load because the two bridge legs must be out of phase, the ZVS condition at no load and high line from Eqs. (8) and (11) is

$$L_M \cdot \left( \frac{V_{IN(max)}}{4 \cdot (2L_M + L_R) \cdot f_S} \right)^2 \geq 4CV_{IN(max)}^2. \quad (12)$$

Finally, from Eq. (12), the approximate value of  $L_M$  required to maintain ZVS at no load and high line is

$$L_M \leq \frac{1}{256Cf_S^2}, \quad (13)$$

assuming that  $2L_M \gg L_R$ .

As can be seen from Fig. 4, current  $i_M$  flowing through magnetizing inductance  $L_M$  introduces a current asymmetry in the two bridge legs. Namely, the coupling of windings DF and FE causes current  $i_2$  flowing through winding FE to equal current  $i_3$  flowing through winding DF, so that  $i_1=i_2+i_M$ . Therefore, in the proposed circuit, leading leg  $S_1$ - $S_2$  always carries a higher current than lagging leg  $S_3$ - $S_4$ ; the difference being magnetizing current  $i_M$ .

Since, in the circuit in Fig. 3, primary current  $i_P=i_1+i_2$  is divided between the two bridge legs, *i.e.*, it is effectively carried by two parallel current paths, the minimum total conduction loss for the primary switches is obtained if equal currents flow through both legs of the bridge. With equal

current distribution between the bridge legs, the total conduction loss of the bridge can be cut in half because the two conducting switches of the bridge carry only one half of primary current  $i_P$ . Therefore, to simultaneously achieve ZVS at no load and minimize the bridge conduction loss in the circuit in Fig. 3, it is necessary to select the maximum magnetizing inductance  $L_M$  determined by Eq. (13) so that the imbalance of the leg currents is minimized. Furthermore, if this minimized current imbalance is still significant, *i.e.*, if current  $i_2$  in lagging leg  $S_3$ - $S_4$  is significantly lower than current  $i_1$  in the leading leg  $S_1$ - $S_2$ , different size switches can be selected for the two legs, which may reduce the cost of the implementation without sacrificing the circuit performance. Phase-shift control of the proposed circuit can be implemented with any of the commercially available phase-shift IC controllers.

### III. EXPERIMENTAL RESULTS

The performance of the proposed inverter circuit was verified on a 1-MHz, 230-W prototype circuit that was designed to operate from a 320 V - 400 V dc input and deliver 27-V<sub>RMS</sub> ac output voltage. The experimental circuit was implemented with the following components: switches  $Q_1$ - $Q_4 =$  IRF820, blocking capacitors  $C_{B1}$  and  $C_{B2} = 0.1 \mu F/400$  V polypropylene capacitor, and resonant capacitor  $C_R = 56$  nF multi-layer ceramic capacitor. The core of transformer TR is a combination of E22/6/16-3F4 and PLT22/16/2.5-3F4. The primary and secondary windings of transformer TR are five turns of Litz wire (150 strands, AWG #42) and a single turn of two parallel connected copper foils (width = 5mm, thickness = 7 mils), respectively. As shown in Fig. 7, resonant inductors  $L_{R1}$  and  $L_{R2}$  and coupled inductor  $L_C$  are integrated on the same core using two

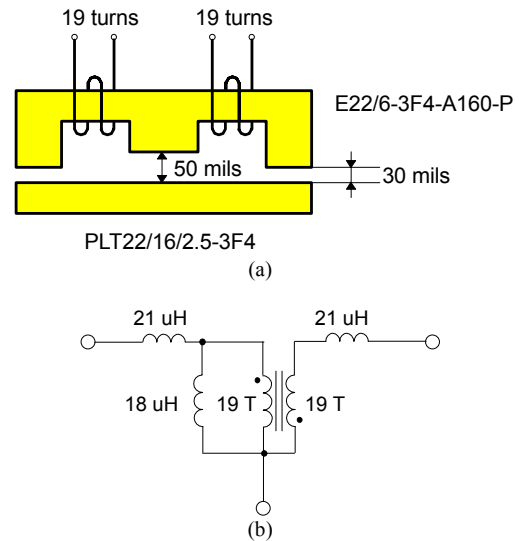


Fig. 7. Integrated magnetic which contains coupled inductor  $L_C$  and resonant inductor  $L_{R1}$  and  $L_{R2}$ : (a) winding structure; (b) equivalent circuit diagram.

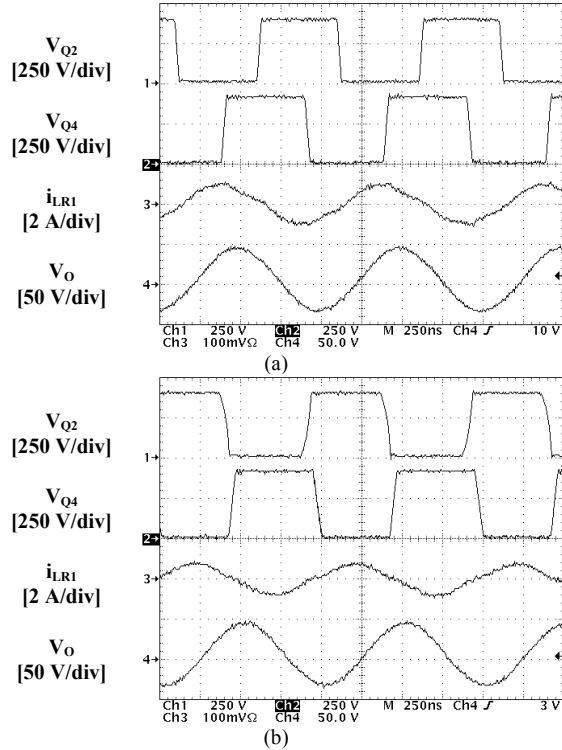


Fig. 8. Measured waveforms at  $V_{IN}=400$  V,  $V_O=27.1$  V<sub>RMS</sub>: (a)  $P_O=230$  W and (b)  $P_O=9.2$  W. Time base: 250 ns/div.

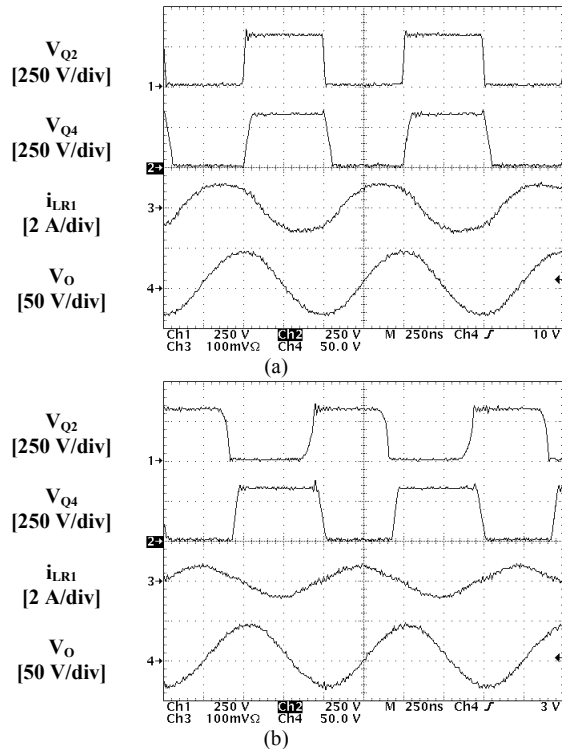


Fig. 9. Measured waveforms at  $V_{IN}=320$  V,  $V_O=27.6$  V<sub>RMS</sub>: (a)  $P_O=230$  W and (b)  $P_O=9.2$  W. Time base: 250 ns/div.

windings with 19 turns each. In fact, the leakage inductance of the coupled inductor in Fig. 7 is utilized as resonant

inductors  $L_{R1}$  and  $L_{R2}$ . The leakage inductance of the coupled inductor is adjusted so that  $L_{R1}=L_{R2} = 21$   $\mu$ H, whereas its magnetizing inductance  $L_M$  is approximately 18  $\mu$ H. The control circuit was implemented with a constant-frequency phase-shift controller UC3875.

Figures 8 and 9 show the measured waveforms of the proposed inverter at full load and 4% load with 320 V and 400 V input voltage, respectively. All the switches operate with ZVS in all conditions. Moreover, the peak-to-peak current of the resonant inductor at 4% load is approximately 65% of that at the full load. As a result, the conduction loss at light load is approximately 42% of that at the full load. It should be noted that the peak-to-peak current of the resonant inductor in the conventional inverter at light load is often greater than that of the full load condition [8].

#### IV. SUMMARY

A new constant-frequency dc-ac resonant inverter that employs a coupled inductor to achieve ZVS in the entire load and input voltage range with a reduced circulating current has been described. The performance of the proposed approach was evaluated on a 230-W inverter that delivers 1-MHz, 27-VRMS ac output voltage from a 320 V - 400 V dc input. The results shows all the switches operate with ZVS in all conditions. Moreover, the conduction losses at light loads are greatly reduced.

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